

# **$M_{wpd}$ : A Duration-Amplitude Procedure for Rapid Determination of Earthquake Magnitude and Tsunamigenic Potential from $P$ Waveforms**

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*Accepted 2008 September 13. Received 2008 September 11; in original form 2007 October 03.*

*Abbreviated title:  $M_{wpd}$ : Duration-Amplitude Magnitude*

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## **Summary**

We present a duration-amplitude procedure for rapid determination of a moment magnitude,  $M_{wpd}$ , for large earthquakes using  $P$ -wave recordings at teleseismic distances.  $M_{wpd}$  can be obtained within 20 minutes or less after the event origin time as the required data is currently available in near-real time. The procedure determines apparent source durations,  $T_0$ , from high-frequency,  $P$ -wave records, and estimates moments through integration of broadband displacement waveforms over the interval  $t_P$  to  $t_P + T_0$ , where  $t_P$  is the  $P$  arrival time. We apply the duration-amplitude methodology to 79 recent, large earthquakes (Global Centroid-Moment Tensor magnitude,  $M_w^{CMT}$ , 6.6 to 9.3) with diverse source types. The results show that a scaling of the moment estimates for interplate thrust and possibly tsunami earthquakes is necessary to best match  $M_w^{CMT}$ . With this scaling,  $M_{wpd}$  matches  $M_w^{CMT}$  typically within  $\pm 0.2$  magnitude units, with a standard deviation of  $\sigma = 0.11$ , equaling or outperforming other approaches to rapid magnitude determination. Furthermore,  $M_{wpd}$  does not exhibit saturation; that is, for the largest events,  $M_{wpd}$  does not systematically underestimate  $M_w^{CMT}$ . The obtained durations and duration-amplitude moments allow rapid estimation of an energy-to-moment parameter  $\Theta^*$  used for identification of tsunami earthquakes. Our results show that  $\Theta^* \leq -5.7$  is an appropriate cutoff for this identification, but also show that neither  $\Theta^*$  nor  $M_w$  is a good indicator for tsunamigenic events in general. For these events we find that a reliable indicator is simply that the duration  $T_0$  is greater than about 50 sec. The explicit use of the source duration for integration of displacement seismograms, the moment scaling, and other characteristics of the duration-amplitude methodology make it an extension of the widely used,  $M_{wp}$ , rapid-magnitude procedure. The need for a moment scaling for interplate thrust and possibly tsunami earthquakes may have important implications for the source physics of these events.

**Key words:** earthquakes, Richter magnitude, seismic moment, seismograms, tsunami, earthquake-source mechanism.

## Introduction

Effective tsunami warning and emergency response for large earthquakes requires accurate knowledge of the event size within 30 minutes or less after the event origin time (OT). The 26 December 2004, M9 Sumatra-Andaman earthquake caused a tsunami that devastated coasts around the Eastern Indian Ocean within 3 hours; the 17 July 2006, M7.7 Java earthquake caused an unexpectedly large and destructive tsunami. For both events the magnitudes available within the first hour after the event origin time severely underestimated the event size (Kerr, 2005; PTWC, 2006ab).

Currently, the earliest, accurate estimates of the size of major and great earthquakes come from long-period, moment-tensor determinations, including the authoritative, Global Centroid-Moment Tensor (CMT) determination and corresponding moment-magnitude,  $M_w^{CMT}$  (Dziewonski *et al.*, 1981; Ekström, 1994), and related procedures (*e.g.*, Kawakatsu, 1995). These estimates are based on seismic  $S$  and surface-wave waveform recordings, typically not available until an hour or more after OT. Other, more rapid moment-tensor based estimates tend to underestimate the size of great earthquakes, as we discuss below.

Another procedure based on surface waves is the mantle magnitude,  $M_m$ , (Okal and Talandier, 1989; Newman and Okal, 1998; Weinstein and Okal, 2005). The spectral amplitude of mantle Rayleigh waves at variable periods (between 50 and 300 sec for large events), combined with approximate corrections for geometrical spreading and Rayleigh wave excitation at the source, gives the  $M_m$  estimate and a corresponding moment.  $M_m$  is potentially available within minutes after the first Rayleigh wave passage (*i.e.* about 20 min after OT at 30° great-circle distance (GCD), and about 50 min after OT at 90° GCD), but for very large events the analysis of waves at increased periods (450 sec or more) may be required (Weinstein and Okal, 2005; UNESCO, 2005) leading to an increased delay after OT for obtaining the  $M_m$  estimate.

Seismic  $P$ -waves are the first signals to arrive at seismic recording stations. At teleseismic distances (30-90° GCD) the arrival times of the initial  $P$ -wave are used routinely to locate the earthquake hypocentre within about 10 to 15 minutes after OT. The initial  $P$ -waves and following  $P$ -wave train also contain comprehensive information about the event size and source character. Boatwright and Choy (1986) show that the total radiated seismic energy can be estimated from the  $P$ -waves alone.

There are a number of procedures for rapid analysis of large earthquakes using seismic  $P$ -waves currently in use at earthquake and tsunami monitoring centers. Because these procedures use only the  $P$ -wave portion of a seismogram, event size estimates are potentially available only a few minutes after the  $P$  waveform has been recorded at teleseismic distances, *i.e.* in as little as 10-15 min after OT at 30° GCD, and about 20 min after OT at 90° GCD.

One of these procedures is the U.S. Geological Survey National Earthquake Information Center (NEIC) Fast Moment Tensor (Sipkin, 1994; <http://earthquake.usgs.gov>) which produces an estimate of the seismic moment tensor and moment magnitude,  $M_w^{NEIC}$ , for earthquakes of magnitude of 5.5 or greater within the order of 30 min after OT through automated processing and inversion of  $P$ -wave waveforms.

Another  $P$ -wave procedure is the widely used,  $M_{wp}$  moment-magnitude algorithm (Tsuboi *et al.*, 1995; Tsuboi *et al.*, 1999; Tsuboi, 2000) which considers very-broadband,  $P$ -wave displacement seismograms as approximate far-field, source-time functions. These displacement seismograms are integrated and corrected approximately for geometrical

spreading and an average radiation pattern to obtain scalar moments at each station. Application of the standard moment magnitude formula, averaging over stations and optionally applying a magnitude dependent correction (Whitmore *et al.*, 2002) gives a moment magnitude,  $M_{wp}$ , for an event.

$M_w^{NEIC}$  and  $M_{wp}$  match closely  $M_w^{CMT}$  up to  $M_w^{CMT} \approx 7.5$ , but at greater magnitudes they tend to increasingly underestimate  $M_w^{CMT}$  (Figure 1, Table 1). To resolve this magnitude saturation problem while providing accurate and rapid magnitude estimates for large earthquakes, a number of authors have proposed new methodologies for magnitude determination based on *P*-wave signals.

Menke and Levin (2005) propose that the ratio of long-period, *P*-wave displacement amplitudes between a target event and a nearby reference event of known size can rapidly provide the magnitude of the target event. Lockwood and Kanamori (2006) show that wavelet analysis of *P*-waves distinguishes a significantly greater amplitude of the long-period, *W*-phase for the 26 December 2004, M9 Sumatra-Andaman relative to the *W*-phase of the 28 March 2005, M8.6 Northern Sumatra earthquakes. (The *W*-phase is a superposition of Rayleigh wave overtones that arrive before the *S* wave.) They propose that such analysis can be used for rapid identification of the largest, great earthquakes and their high tsunamigenic potential.

Bormann and Wylegalla (2005), Bormann *et al.* (2006) and Bormann and Saul (2008) calculate a cumulative  $m_B$  magnitude,  $mBc$ , by summing the peak velocity amplitudes for major signal pulses between consecutive zero crossings in the *P* waveform. Hara (2007) combines measures of the high-frequency duration and maximum displacement amplitude of *P*-waveforms for a set of large, shallow earthquakes to determine an empirical relation for moment magnitude.

Lomax (2005) shows for very large earthquakes that the location of the end of rupture, and thus an estimate of the event size, can be rapidly determined from measures of the *P*-wave duration on high-frequency records. Lomax and Michelini (2005) note that the ratio of the high-frequency, *P*-wave durations from the 2004, M9 Sumatra-Andaman and the 2005, M8.6 Northern Sumatra earthquakes match the ratio of the CMT moment values for the two events, and suggest that the high-frequency, *P*-wave duration could be used for rapid magnitude estimation for individual events. Lomax *et al.* (2007) use teleseismic ( $GCD \geq 30^\circ$ ), *P*-wave signals to estimate radiated seismic energy,  $E$ , and source duration,  $T_0$ , and show that an *energy-duration* moment relation,  $M_0^{ED} \propto E^{1/2} T_0^{3/2}$ , based on an expression for  $E$  from Vassiliou and Kanamori (1982), gives a moment magnitude,  $M_{ED}$ , that matches closely  $M_w^{CMT}$  for a set of recent, large earthquakes.

These new methodologies for rapid magnitude determination based on *P*-wave signals all produce useful magnitude estimates,  $M^{est}$ , for very large earthquakes. Most of these methodologies, however, show significant differences with  $M_w^{CMT}$  (*i.e.*,  $|M^{est} - M_w^{CMT}| \geq 0.3$ ) for many events, including some of the most important and destructive interplate thrust events and tsunami earthquakes (tsunami earthquakes are characterized by unusually large tsunamis and a deficiency in moment release at high frequencies, *e.g.*, Kanamori, 1972; Polet and Kanamori, 2000; Satake, 2002). Most of these methodologies also give  $\Delta M = M^{est} - M_w^{CMT}$  values which change systematically with increasing  $M_w^{CMT}$ ; this effect is equivalent to the magnitude saturation of  $M_w^{NEIC}$  and  $M_{wp}$ .

To further investigate and resolve these problems, we introduce here a rapid and robust,

*duration-amplitude* procedure to obtain an earthquake moment and a moment magnitude,  $M_{wpd}$ , from  $P$ -wave recordings at teleseismic distances. This procedure first determines apparent source durations,  $T_0$ , from high-frequency,  $P$ -wave records, and then estimates moments through integration of broadband displacement records over the interval  $t_P$  to  $t_P+T_0$ , where  $t_P$  is the  $P$  arrival time. This methodology can be viewed as an extension of the  $M_{wp}$  moment-magnitude algorithm.

We begin by presenting the theoretical development and calibration of the duration-amplitude procedure using a set of recent, large earthquakes with diverse source types and ideal knowledge of their source parameters (*e.g.*, depth, tectonic setting and mechanism). Next we discuss practical application of the procedure with regards to rapid determination of the source parameters. Finally we examine the performance of the duration-amplitude procedure and  $M_{wpd}$  magnitudes, and the related estimates of tsunamigenic potential, and we present hypothesis on physical implications of this procedure for large interplate thrust earthquakes.

## Theoretical development and calibration of the duration-amplitude procedure

### Basic theory

Given the far-field,  $P$ -displacement  $u(t)$ , for an earthquake source of rupture duration,  $T_0$ , a well established theoretical expression for the scalar, seismic moment,  $M_0$ , is,

$$M_0 = C_M \int_{t_P}^{t_P+T_0} u(t) dt, \quad (1)$$

where  $t_P$  is the  $P$  arrival time,  $u(t)$  is corrected for geometrical spreading and attenuation, and  $C_M$  is a constant that depends on the density and wave speed at the source and station, a double-couple radiation pattern and other factors (*e.g.*, Aki and Richards, 1980; Boatwright and Choy, 1986; Tsuboi *et al.*, 1995; Newman and Okal, 1998; Kanamori and Rivera, 2004; see Appendix A for details). Equation (1) suggests that the scalar moment,  $M_0$ , of an earthquake can be determined from  $P$ -wave, displacement seismograms. Application of the standard moment-magnitude formula to the obtained  $M_0$ ,

$$M_w = (\log_{10} M_0 - 9.1) / 1.5, \quad (2)$$

(Hanks and Kanamori, 1979; Bormann, 2002) gives a  $P$ -wave estimate of the moment magnitude,  $M_w$ , for an event.

Equation (1) cannot be used directly to obtain accurate moment estimates for a number of reasons, including the presence of surface reflected and other secondary phases, and the difficulty of estimating  $T_0$ . The  $M_{wp}$  magnitude procedure addresses some of these problems by estimating the scalar moment from the larger of the first peak or the first peak-to-peak amplitudes on  $P$ -displacement seismograms integrated using Equation (1), though the integral is performed without explicit knowledge or use of  $T_0$  (*e.g.*, Tsuboi *et al.*, 1999).

To make further use of Equation (1) to obtain more accurate, rapid moment-magnitude estimates, we begin by examining moments,  $\hat{M}_0$ , and magnitudes,  $M_{wpd}$ , determined through application to teleseismic  $P$ -wave, ground-displacement seismograms of a modified form of Equation (1),

$$\hat{M}_0 = k C_M \text{Max} \left[ \int_{t_p}^{t_p+T_0} u^+(t) dt, \int_{t_p}^{t_p+T_0} |u^-(t)| dt \right]. \quad (3)$$

The modifications in Equation (3) includes the following: 1) The integral in Equation (1) is taken separately over the positive,  $u^+(t)$ , and the absolute value of negative,  $|u^-(t)|$ , displacement amplitudes to help separate the direct  $P$  waves from surface reflection phases and other phases with opposite polarity; the maximum of these two integrals is used to calculate the moment estimate. 2) A constant,  $k$ , is included to compensate for unknown errors and biases in the terms of  $C_M$  and in the correction of  $u(t)$  for attenuation and geometrical spreading (if  $C_M$  and the corrections were physically exact, a value of  $k = 1$  would be expected). In addition, the source duration,  $T_0$ , is estimated through measures on high-frequency,  $P$ -wave seismograms (Lomax, 2005; Lomax *et al.*, 2007) and explicitly used to define the upper limit of integration. Application of the standard moment-magnitude formula, Equation (2), using  $\hat{M}_0$  and averaging over stations using robust statistics (20% trimmed mean) gives a  $P$ -wave moment magnitude,  $M_{wpd}$ , for an event.

Further details on this procedure are given in Appendices A, B and C; the processing steps are illustrated in Figure 3. We note here that the amplitude correction of the displacement waveforms for attenuation and geometrical spreading and the calculation of  $C_M$  make use of the PREM model (Dziewonski and Anderson, 1981) without a crust (hereinafter referred to as PREM\_NC), since most large events occur in oceanic regions. For shallow continental events, the effect of the crust on  $C_M$  is introduced as a magnitude correction using the PREM properties for the lower crust. Also, the radiation pattern factor in  $C_M$  for strike-slip events, which differs greatly from that for all other event types, is determined empirically. Table 1 indicates the classification of each event according to source type and oceanic versus continental setting, mainly based on event information from the NEIC (<http://earthquake.usgs.gov>) and the Global CMT Catalog (<http://www.globalcmt.org>). This classification takes into account the epicenter, depth and moment-tensor mechanism in relation to the background seismicity and the surrounding tectonic plates and plate boundaries; in a few cases additional information from the NEIC tectonic summary is used.

Our use in Equation (3) of the maximum of the integrals over positive and negative  $P$ -displacement is a direct extension to all peaks in the interval  $T_0$  after  $P$  of the use in the  $M_{wp}$  procedure of the first peak or the first peak-to-peak of the displacement integral (*e.g.*, Tsuboi *et al.*, 1999). It is difficult, if not impossible, to justify this procedure theoretically for all event types, event depths and  $P$ -group phases. However, we find that the use of this procedure, relative to integrating the absolute value of the displacement, gives better agreement with  $M_w^{CMT}$  magnitudes, and a value of the constant  $k$  in Equation (3) that is closer to the ideal value of 1.

### ***Direct application to large earthquakes***

Figure 4 shows a comparison of the obtained magnitudes,  $M_{wpd}$ , with  $M_w^{CMT}$  for 79, recent, large earthquakes ( $M_w^{CMT}$  6.6 to 9.3; Figure 2 and Table 1) using no knowledge of the event type (Figure 4a) and using ideal knowledge of the depth, tectonic setting and mechanism for each event (Figures 4b and 4c). This comparison shows that  $M_{wpd}$  matches closely  $M_w^{CMT}$  up to  $M_w^{CMT} \sim 7.5$ , but with increasing magnitude  $M_{wpd}$  tends to increasingly underestimate  $M_w^{CMT}$ . This is a similar result as obtained for  $M_{wp}$  (Figure 1), though  $M_{wp}$  gives an even larger underestimate than  $M_{wpd}$  of  $M_w^{CMT}$  above  $M_w^{CMT} \sim 7.5$ , primarily because  $M_{wp}$  only considers the first part of the  $P$  wave train while  $M_{wpd}$  is based on the full interval of duration  $T_0$  after the  $P$

arrival. The NEIC Fast Moment Tensor magnitude,  $M_w^{NEIC}$ , (Sipkin, 1994; <http://earthquake.usgs.gov>), based on waveform inversion, also shows an increasing underestimate of  $M_w^{CMT}$  above  $M_w^{CMT} \sim 7.5$  (Figure 1).

Closer examination of Figure 4 shows that the trend of increasing underestimate of  $M_w^{CMT}$  by  $M_{wpd}$  (i.e.  $\Delta M = M_{wpd} - M_w^{CMT}$  becomes more negative) with increasing  $M_w^{CMT}$  occurs mainly for *interplate thrust* earthquakes (type I in Table 1).  $M_{wpd}$  matches well  $M_w^{CMT}$  for most events of other types, agreeing over a wide range of magnitudes for strike-slip (types S and So), intraplate (type P), intermediate depth (downdip, type W) and deep earthquakes (type D), and over the limited range of available magnitudes for reverse-faulting (type R and Ro) and normal-faulting (type N and No) crustal earthquakes. It cannot be excluded that tsunami earthquakes (type T) follow a trend similar to that of interplate thrust earthquakes, due to the lack of large tsunami earthquakes.

Thus we find for larger ( $M_w^{CMT} > \sim 7.5$ ) interplate thrust events that the moments determined from the *P*-wave train through application of Equation (3), and apparently also through *P*-waveform inversion (e.g.,  $M_w^{NEIC}$ , Figure 1a), underestimate the corresponding CMT moments, derived from inversion of long period *S* and surfaces wave.

### ***Moment scaling for interplate thrust and tsunami earthquakes***

The variation of  $\Delta M = M_{wpd} - M_w^{CMT}$  differences for interplate thrust earthquakes as a function of  $M_w^{CMT}$  (Figure 4c) and a similar variation as a function of  $M_{wpd}$  suggest that more accurate moment estimates for these events,  $M_0^I$ , can be obtained by scaling  $\hat{M}_0$  with a factor composed of  $\hat{M}_0$  raised to some power, i.e.,

$$M_0^I = \hat{M}_0 \left( \frac{\hat{M}_0}{M_0^{cutoff}} \right)^R, \quad (4)$$

where  $\hat{M}_0$  is given by Equation (3) and  $M_0^{cutoff}$  is a constant cutoff moment below which the scaling is not applied. We also apply the moment scaling, Equation (4), to tsunami earthquakes, since these events fall within the trend of  $\Delta M$  differences for interplate thrust earthquakes and because it is difficult to distinguish these two types of events in near real-time analysis. Application of the standard moment-magnitude formula, Equation (2), and averaging over stations gives the corresponding *P*-wave moment magnitude,  $M_{wpd}$ . (see Appendix B for further details)

### ***Application with moment scaling to large earthquakes***

Application of Equation (4) to the interplate thrust and tsunami events from the set of studied earthquakes over a range of values of  $R$  and  $M_0^{cutoff}$  gives  $R \approx 0.45$  and  $M_0^{cutoff} \approx 7.5 \times 10^{19}$  N-m (equivalent to  $M_w \approx 7.2$ ) for the best match of  $M_{wpd}$  to  $M_w^{CMT}$ . (The optimal value of  $M_0^{cutoff}$  and  $R$  are sensitive to the algorithms used to estimate  $T_0$  and moment, see Appendix B). Thus we arrive at a preferred, duration-amplitude expression for moment estimation,

$$M_0^{pd} = \hat{M}_0 \left( \frac{\hat{M}_0}{M_0^{cutoff}} \right)^{0.45}, \quad (5a)$$

for interplate thrust and tsunami events with  $\hat{M}_0 \geq M_0^{cutoff}$ , and

$$M_0^{pd} = \hat{M}_0, \quad (5b)$$

otherwise, where  $\hat{M}_0$  is given by Equation (3) with  $C_M=1.62 \times 10^{19}$  and  $k \approx 1.2$  (see Appendices A and B; see Table 2 for depth corrections).  $M_{wpd}$  magnitudes determined using Equations (5a), (5b) and (2) for the studied earthquakes are shown in Figure 5 and Table 1. These results show that  $M_{wpd}$ , with moment scaling for interplate thrust and tsunami events, matches  $M_w^{CMT}$  typically within  $\pm 0.2$  magnitude units, with a standard deviation of only  $\sigma=0.11$ .

### ***Estimation of energy-to-moment parameter $\Theta^*$***

The energy-to-moment parameter,  $\Theta$ , (e.g., Newman, and Okal, 1998; Weinstein and Okal, 2005) for identification of tsunami earthquakes is defined as,

$$\Theta = \log_{10} \frac{E}{M_0}, \quad (6)$$

where  $E$  is the radiated seismic energy and  $M_0$  the moment. Weinstein and Okal (2005) note that standard earthquake scaling laws (assuming a constant stress drop) predict a value of  $\Theta \approx -4.9$ , but find  $\Theta$  values around  $-6.0$  or less for tsunami earthquakes. Thus anomalously low values of a rapid estimate of  $\Theta$ , combined with knowledge of an earthquake's location, size, tectonic setting and likely source type, can be an important indicator of a potential tsunami earthquake.

From duration-amplitude estimates of moment,  $M_0^{pd}$ , and duration,  $T_0$ , we can obtain an approximation to  $\Theta$ ,  $\Theta^*$ , through application of the energy-duration relation of Lomax et al. (2007),

$$M_0^{ED} = c E^{1/2} T_0^{3/2}, \quad (7)$$

where  $c \approx 1.55 \times 10^{10}$  for average crust - upper mantle material properties. Substituting  $M_0^{pd}$  for  $M_0^{ED}$  in Equation 7, solving for  $E$  and substituting into Equation 6, gives,

$$\Theta^* = \log_{10} (c^{-2} M_0^{pd} / T_0^3). \quad (8)$$

The approximate  $\Theta^*$  values should only be used when the uncertainty  $\sigma_{T_0}$  in  $T_0$  is small, since the dependence of  $\Theta^*$  on  $T_0^{-3}$  amplifies error in  $T_0$  into  $\Theta^*$ .  $\Theta^*$  values determined using Equation (8) for the studied earthquakes where  $\sigma_{T_0} < 2T_0/3$  are listed in Table 1 and plotted in Figure 6 as function of  $M_{wpd}$ .

## **Practical application of the duration-amplitude procedure**

Without moment scaling (Equation (5a))  $M_{wpd}$  provides a closer match to  $M_w^{CMT}$  magnitude, including for larger interplate thrust events and tsunami earthquakes, than do other procedures for rapid magnitude estimation (standard deviation of  $\sigma=0.17$ ; cf. Figures 1, 4b and 4c, Table 1). Furthermore, a “raw”  $M_{wpd}$  given by direct application of Equation (3) without any corrections for event type (e.g., no crustal correction for shallow continental events, no correction for radiation pattern for strike-slip events) still matches  $M_w^{CMT}$  with  $\sigma=0.18$  (Figure 4a). However, as with all rapid analysis methodologies based on body-wave signals, knowledge of the hypocenter location, the tectonic setting and likely focal mechanism is needed to obtain the best and most informative results.

## Rapid identification of event type and other source parameters

Obtaining the best match of  $M_{wpd}$  to  $M_w^{CMT}$  requires identification of interplate thrust and tsunami earthquakes for application of the moment scaling, and a reliable depth estimate and further classification by type, *e.g.*, as continental, oceanic, strike-slip, or deep, for application of corrections for PREM properties at the source depth (Table 2). The correction for PREM properties at the source depth relative to average, upper-mantle properties is only significant (*i.e.*, gives a magnitude change  $\delta M > 0.1$ ) for events deeper than 400 km and for continental crustal events. Also, the corrections for continental crustal type ( $\delta M \approx -0.15$ ) and for strike-slip mechanism ( $\delta M \approx 0.13$ ) approximately cancel for continental, strike-slip events. However, not all events can be easily classified within minutes after OT. For example, we classify the 12 September 2007, 23:49, M7.9 Indonesia earthquake as a downdip event (based on the epicentral location and the CMT centroid depth of 44 km; Table 1) giving  $M_{wpd} = 7.9$ ; but the epicentral location and shallow initial depth estimate for this event could imply that it is an interplate thrust event, in which case the amplitude-duration moment scaling should be applied, giving  $M_{wpd} = 8.2$ .

In the near future, information on the hypocenter location, tectonic setting and likely focal mechanism of an event should be available before the duration-amplitude analysis is performed, thus likely interplate thrust and tsunami events, the event type and the approximate source depth can be identified rapidly. Currently, the epicenter for most events can be determined accurately within minutes of OT; the main difficulties lie with the determination of the hypocentral depth and, secondarily, of the source mechanism. Improvements in depth determination based on standard earthquake location procedures are not likely, due to fundamental limitations of the ray coverage at the source of the rapidly available, first *P* arrival data. Instead, improved depth estimation may come from prior information on the depth of seismicity and plate boundaries (*e.g.*, Hayes and Wald, 2008) which can provide a likelihood function for depth based on the epicentral location. Similarly, maps of crustal types (*e.g.*, Mooney *et al.*, 1998; Bassin *et al.*, 2000) can provide a likelihood function for tectonic setting based on the epicentral location. For the determination of  $M_{wpd}$  knowledge of the source mechanism, though less important than event depth, could provide further constraint on the tectonic setting and event type (*e.g.*, distinguishing between interplate thrust and normal-faulting, outer-rise earthquakes, both of which occur near subduction zones). Rapid and robust estimation of mechanism may be possible using existing procedures based on the first-motions and initial amplitudes of *P*-waveforms.

## Discussion

We have introduced a duration-amplitude procedure to obtain rapidly an earthquake moment,  $M_0^{pd}$ , and moment magnitude,  $M_{wpd}$ , from *P*-wave recordings at teleseismic distances. Because the required recordings are available in near-real time at earthquake and tsunami monitoring centers,  $M_{wpd}$  can be available within about 20 minutes after OT. For major and great earthquakes ( $M_w^{CMT} \geq 7.0$ ),  $M_{wpd}$  (with moment scaling for interplate thrust and tsunami events) matches  $M_w^{CMT}$  typically within  $\pm 0.2$  magnitude units, with a standard deviation of only  $\sigma = 0.11$  (Figure 5, Table 1). In addition,  $M_{wpd}$  does not exhibit saturation; that is, for the largest events,  $M_{wpd}$  does not systematically underestimate  $M_w^{CMT}$  and  $\Delta M = M_{wpd} - M_w^{CMT}$  remains small. Thus  $M_{wpd}$  equals or outperforms other procedures for rapid moment magnitude determination. The results of other procedures, using different and smaller sets of events than used here, are:

- $M_{ED}$  (Lomax *et al.*, 2007) matches  $M_w^{CMT}$  typically within  $\pm 0.3$  magnitude units, with  $\sigma = 0.16$ , and  $\Delta M = M_{ED}^{est} - M_w^{CMT}$  for  $M_{ED}$  does not change with increasing  $M_w^{CMT}$ ;



- *mBc* (Bormann *et al.*, 2006; Bormann and Saul, 2008) matches  $M_w^{CMT}$  typically within  $\pm 0.5$  magnitude units, with  $\sigma=0.26$ , but there is a trend in  $\Delta M$  for uncorrected *mBc* to become more positive with decreasing  $M_w^{CMT}$  (the large  $\sigma$  with respect to  $M_w^{CMT}$  is expected since *mBc* is fundamentally an energy magnitude, while the trend in  $\Delta M$  can be compensated via a regression relation given by Bormann and Saul, 2008);
- the rapid magnitude estimates of Hara (2007) shows a match with  $M_w^{CMT}$  typically within  $\pm 0.3$  magnitude units, with  $\sigma=0.18$ , and  $\Delta M$  for this magnitude is stable or possibly becomes more negative with increasing  $M_w^{CMT}$ ;
- our corrected  $M_{wp}$  results (Figure 1c; Table 1) match  $M_w^{CMT}$  typically within  $\pm 0.5$  magnitude units, with  $\sigma=0.25$ , and  $\Delta M$  for  $M_{wp}$  becomes rapidly increasingly negative with increasing  $M_w^{CMT}$ .

The improved agreement between  $M_{wpd}$  and  $M_w^{CMT}$  relative to other rapid procedures, including  $M_{wp}$ , can be attributed primarily to the use in Equation (3) of the full  $t_P$  to  $t_P+T_0$  interval for integration with testing of integrals over positive and negative values of displacement, and to the application of the moment scaling, Equation (5a), for interplate thrust and tsunami earthquakes. This agreement is also dependent on the use of additional corrections for certain events types, and a robust procedure for estimating  $T_0$  from high-frequency seismograms (see Appendices A and B for details). Indeed, much of the scatter in  $M_{wpd}$  versus  $M_w^{CMT}$  for  $M_w < \sim 7.2$  (Figures 4 and 5) can be attributed to large errors in the  $T_0$  estimation. The  $M_{wpd}$  results indicate that testing of the integral in Equation (3) over positive and negative values of displacement separates adequately the direct *P* waves from surface reflection phases and other secondary phases, even when the rupture duration,  $T_0$ , is large. This testing is analogous to the selection in the  $M_{wp}$  magnitude procedure of the larger of the first peak or the first peak-to-peak amplitude of the integral Equation (1). The moment scaling used here is likely related to the magnitude dependent correction to  $M_{wp}$  proposed by Whitmore *et al.* (2002) and to the values of the coefficients in the regression of Hara (2007), in both cases applied to all earthquakes. In contrast, we find here that moment scaling is only needed for interplate thrust earthquakes, and possibly for tsunami earthquakes. The characteristics of the duration-amplitude procedure noted above show that it is an extension of the  $M_{wp}$  moment-magnitude algorithm, recalling also that both procedures are ultimately based on Equation (1).

### ***Energy-to-moment parameter $\Theta^*$ and tsunamigenic potential***

We have shown that the duration-amplitude estimates of moment,  $M_0^{pd}$ , and duration,  $T_0$ , can be combined with the energy-duration relation of Lomax *et al.* (2007) to provide a rapid approximation,  $\Theta^*$ , (Equation 8) to the energy-to-moment parameter  $\Theta$  used for identification of tsunami earthquakes (*e.g.*, Newman, and Okal, 1998; Weinstein and Okal, 2005). Duration-amplitude estimates of  $\Theta^*$  using Equation 8 are listed in Table 1 and are shown in Figure 6. To simulate the results that would be available with rapid application of the duration-amplitude procedure, we show  $\Theta^*$  values only for events where  $\sigma_{T_0} < 2T_0/3$  (a stronger cutoff would be used in practice) and plot  $\Theta^*$  against  $M_{wpd}$  and not  $M_w^{CMT}$ , which is not available rapidly.

The duration-amplitude estimates of  $\Theta^*$  are  $\Theta^* \leq -5.8$  for all studied tsunami earthquakes, thus  $\Theta^* \leq -5.7$  may be an appropriate cutoff for identification of these events (Figure 6). Some interplate thrust, downdip and strike-slip events have low  $\Theta^*$  values ( $\Theta^* \leq -5.5$ ), and deep events have high  $\Theta^*$  values.  $\Theta^*$  is low,  $\Theta^* = -6.1$ , for a tsunamigenic, interplate thrust event (1998.07.17 Papua New Guinea) that is considered not to be a tsunami earthquake (Heinrich

*et al.*, 2001; Okal, 2003). Low values of  $\Theta^*$  for strike-slip earthquakes can be attributed to overestimate of  $T_0$  for smaller events, perhaps related to the strike-slip radiation pattern producing anomalously low amplitudes and an excessively long coda in the high frequency seismograms used to estimate  $T_0$ . Weinstein and Okal (2005) also find anomalously low values of  $\Theta$  for several strike-slip events. Similarly, the low value  $\Theta^* = -6.3$  for a down-dip earthquake (W; 2005.09.09 New Ireland) can be attributed to overestimate of  $T_0$  for this event due to anomalously high-frequency signal in the depth phases  $pP$  and  $sP$ .

We include in Figure 6 an approximate measure of tsunami importance,  $I_t$  based on maximum water height in meters,  $h$ , and 0-4 descriptive indices,  $i$ , of tsunami effects (deaths, injuries, damage, houses destroyed) from the NOAA/WDC Historical Tsunami Database (NGDC, 2008),

$$I_t = h + i_{deaths} + i_{injuries} + i_{damage} + i_{houses\ destroyed} . \quad (9)$$

For completeness, two events with  $\sigma_{T_0} > 2T_0/3$  but which have  $I_t \geq 1$  are also indicated in Figure 6 (1999.08.17,  $M_w$ 7.6, Turkey,  $I_t=8$ ; 2003.01.22,  $M_w$ 7.5, Mexico,  $I_t=1$ ). Strikingly, Figure 6 shows no clear relationship in between  $I_t$  and  $\Theta^*$ , indicating that  $\Theta$ , while a robust indicator for tsunami earthquakes, is not a good indicator for tsunamigenic potential in general. Similarly, there is no clear relationship between  $I_t$  and  $M_w$ , as represented by  $M_{wpd}$ . Instead, we find that the majority of tsunamigenic events fall to one side of diagonal lines of constant  $T_0$ , as defined by Equation (8). A good separation between events with  $I_t < 1$  (unlikely tsunamigenic events) and those with  $I_t > 1$  (possible tsunamigenic events) is given by the line  $T_0 = 50$  sec. The only major exception is 2003.05.21,  $M_w$ 6.8, N Algeria, a submarine, shallow, thrust event that produced larger than expected tsunami waves, perhaps due to focussing of tsunami waves or slope failure near the source (Hébert and Alasset, 2003). These results suggest that a value of  $T_0 \geq 50$  sec obtained with the rapid, duration-amplitude procedure, if the uncertainty in  $T_0$  is low, is a reliable indicator of a possible destructive, tsunamigenic event.

The importance of the  $T_0$  estimate for the determination of the tsunamigenic potential of an earthquake, and to a lesser degree for the determination of  $M_{wpd}$ , combined with the large uncertainties in  $T_0$  obtained for smaller events and certain event types, indicates a need for future work on improving the accuracy and robustness of the  $T_0$  determination.

### ***Application at local and regional distances***

The duration-amplitude methodology may be applicable at local and regional distances, *i.e.*  $GCD < 30^\circ$ , thus reducing the time delay after OT for obtaining size estimates for larger events. However, relative to the teleseismic analysis presented in this paper, there are many complications when working at local and regional distances. The main difficulty is that significant  $S$  signal may remain on the high-frequency,  $P$ -wave seismograms used for determination of the duration,  $T_0$ , which complicates the analysis of larger and longer duration events. In this case, the direct  $P$ -wave radiation can often be isolated by applying the narrow-band, Gaussian filtering at higher frequencies (e.g., 5-20 Hz), but this requires that high dynamic-range, high sample-rate data is available. At regional distances, there may be additional difficulties due to the multitude of direct, reflected, refracted and converted  $P$  and  $S$  wave types that can contribute to the  $P$  wave train.

### ***Physical implications of moment scaling***

The increasing underestimate of  $M_w^{CMT}$  with increasing magnitude by unscaled  $M_{wpd}$  for interplate thrust (and possibly tsunami) earthquakes (Figure 4) and the consequent need for

moment scaling (Equation 5a) may have important physical implications. The increasing underestimate of  $M_w^{CMT}$  is probably not due to station site or path effects, since then it would occur for all event types, and it is probably not a direct effect of the source mechanism radiation pattern, since then it would not vary with event size. In addition, examination of  $M_{wpd}$  estimates obtained with different long-period cutoffs (Appendix C) indicates that the increasing underestimate of  $M_w^{CMT}$  is not due to magnitude saturation due to insufficient, long-period signal. Thus the increasing underestimate of  $M_w^{CMT}$  may be associated with near-source, dynamic phenomena unique to larger interplate thrust (and possibly tsunami) earthquakes, events which occur at shallow depths. The form of the moment scaling, Equation (5a), suggests a deficiency that increases with event size in the amplitude of far-field, radiated  $P$ -waves relative to the amplitudes expected from the CMT results.

The destructive interference of  $pP$  or  $sP$  waves with direct, down-going  $P$  waves is an often cited explanation for reduced, far-field  $P$  amplitudes. This is, however, a kinematic mechanism which, for large, shallow earthquakes, must be cast into a dynamic framework where the interference will occur within the rupture volume and simultaneous with rupture. The deficiency in amplitude may therefore be associated with a near-field mechanism which reduces the radiated kinetic energy while maintaining the seismic energy balance. A candidate mechanism would be excessive dissipation of the strain energy released during faulting by gravitational, fracturing and frictional processes on or near the fault, alimanted by complex wave interactions around the rupturing fault. Such interactions could involve waves reflected, generated or trapped near the free surface, such as the near-field analogues of  $pP$  and  $sP$ , which may interfere destructively with the fault displacements that produce far-field  $P$ -waves, reducing the amplitude of these waves. We can then hypothesises a transfer of kinetic energy along strike and in the direction of rupture (for long thrust faults) by waves from earlier rupture, producing dynamic stress loading across the fault around the rupture front and augmenting the loading due to nearby fault displacements.

Such dynamic loading near the rupture front could raise the shear stress above the failure yield stress (*e.g.*, Scholz, 2002), decrease the normal stress and thus decrease the effective yield stress (*e.g.*, Oglesby *et al.*, 2000), or drive rupture in zones with a velocity-strengthening friction behavior (*e.g.*, Scholz, 1998). In all these cases, increased fracture, rupture and slip would be induced at the rupture front, including on parts of the fault for which the initial shear stress was much less than the static yield stress, or which have velocity-strengthening behavior, likely in the shallower, up-dip parts of subduction thrusts (*e.g.*, Scholz, 1998). Thus the moment scaling could be a manifestation of a “self-driving” mechanism for large interplate thrust (and possibly tsunami) earthquakes in which an anomalously large proportion of the energy released during rupture is re-absorbed locally to further drive the rupture, and thus to make the earthquake large.

## Conclusions

We have presented a duration-amplitude procedure for determination of a moment magnitude,  $M_{wpd}$ , for large earthquakes within 20 minutes of the event origin time using teleseismic  $P$ -wave recordings. We find that a scaling of the moment estimates for interplate thrust and possibly tsunami earthquakes is necessary to best match  $M_w^{CMT}$ . With this scaling,  $M_{wpd}$  equals or outperforms other approaches to rapid magnitude determination, and does not exhibit saturation.

The characteristics of the duration-amplitude methodology make it an extension of the widely used,  $M_{wp}$ , rapid-magnitude procedure. The need for a moment scaling for interplate thrust and possibly tsunami earthquakes may have important implications for the source physics of

these events.

As with other rapid, earthquake-analysis procedures, obtaining the best match of  $M_{wpd}$  to  $M_w^{CMT}$  requires identification of the event type, a reliable depth estimate, and other source parameters. Many groups are currently working on providing this information more reliably and more rapidly.

The duration-amplitude procedure allows rapid estimation of an energy-to-moment parameter  $\Theta^*$  used for identification of tsunami earthquakes. However, our results show that neither  $\Theta^*$  nor  $M_w$  is a good indicator for tsunamigenic events in general. For these events we find that a reliable indicator is simply that the duration-amplitude duration,  $T_0$ , is greater than about 50 sec.

## Acknowledgements

This work benefited greatly from discussions with Peter Bormann, Massimo Cocco, Paul Earle, Goran Ekström, Barry Hirshorn, Chris Marone, Stefan Nielsen, and Martin Vallée, and from thorough reviews by two anonymous reviewers. The work of A.L. was supported by personal funds; A.M. has been supported by the INGV-DPC (Dipartimento della Protezione Civile) S4 project - “Stima dello scuotimento in tempo reale e quasi-reale per terremoti significativi in territorio nazionale”. We use the Java program SeisGram2K (<http://www.alomax.net/software>) for seismogram analysis, processing and figures, and OpenOffice.org Calc for graphs. The IRIS DMC (<http://www.iris.edu>) provided access to waveforms used in this study.

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Table 1 – Events used in this study and duration-amplitude results

Origin time	Event	Type*	NEIC				CMT			this study		this study, duration-amplitude results				
			latitude (°)	longitude (°)	depth (km)	$M_w^{NEIC}$	depth (km)	$T_0^\dagger$ (sec)	$M_w^{CMT}$	$M_{wp}$	$M_{wp}$ corr $\ddagger$	$T_0$ (sec)	$M_{wpd}$	moment scaling	$M_{wpd}$ scaled	$\Theta^*$
1992.09.02 00:15	Nicaragua	T	11.74	-87.34	44	6.7	15	89	7.6	7.3	7.4	222	7.5	Yes	7.6	-6.9
1992.12.12 05:29	Flores Indonesia	P	-8.48	121.90	49	7.4	20	45	7.7	7.7	7.9	100	7.7		7.7	-5.7
1993.07.12 13:17	Hokkaido	P	42.85	139.20	18	7.3	17	50	7.7	7.6	7.8	87	7.8		7.8	-5.4
1994.01.17 12:30	S California	R	34.21	-118.54	21	6.7	17	16	6.6	6.9	6.9	15	6.7		6.7	-4.8
1994.06.02 18:17	Java	T	-10.48	112.84	6	7.7	15	78	7.7	7.5	7.7	108	7.5	Yes	7.7	-5.9
1994.06.09 00:33	Bolivia	D	-13.84	-67.55	631	8.1	647	58	8.2	7.8	8.0	42	8.2		8.2	-3.9
1994.10.04 13:23	Kuril Islands	P	43.77	147.32	61	8.1	68	60	8.3	7.8	8.1	67	8.2		8.2	-4.5
1995.07.30 05:11	Chile	I	-23.34	-70.29	9	7.9	29	67	8.0	7.6	7.8	101	7.8	Yes	8.0	-5.2
1995.10.09 15:35	Mexico	I	19.06	-104.21	4	7.9	15	66	8.0	7.4	7.6	77	7.6	Yes	7.8	-5.2
1995.12.03 18:01	Kuril Islands	I	44.66	149.30	23	7.6	26	57	7.9	7.6	7.8	61	7.6	Yes	7.8	-5.0
1996.02.17 05:59	Irian Jaya	I	-0.89	136.95	11	8.1	15	66	8.2	N/A	N/A	112	7.8	Yes	8.1	-5.3
1996.02.21 12:51	Peru	T	-9.59	-79.59	4	7.4	15	45	7.5	7.3	7.5	101	7.2	Yes	7.3	-6.4
1998.03.25 03:12	Balleny Islands	So	-62.87	149.52	10	7.8	29	75	8.1	7.8	8.1	114	8.0		8.0	-5.4
1998.07.17 08:49	Papua New Guinea	I	-2.96	141.93	7	7.0	15	39	7.0	6.9	6.9	60	7.0		7.0	-6.1
1999.04.08 13:10	Russia-China	D	43.61	130.35	576	7.1	575	12	7.1	7.0	7.0	10	7.1		7.1	-3.6
1999.08.17 00:01	Turkey	S	40.75	29.86	13	7.4	17	22	7.6	7.6	7.7	67	7.4		7.4	-5.6
1999.09.20 17:47	Taiwan	Ro	23.77	120.98	8	7.4	21	34	7.6	7.6	7.8	62	7.6		7.6	-5.3
1999.10.16 09:46	S California	S	34.59	-116.27	20	7.1	15	30	7.1	7.4	7.5	49	7.0		7.0	
2000.06.04 16:28	Sumatra	P	-4.72	102.09	7	7.7	44	41	7.8	7.8	8.1	87	7.9		7.9	-5.3
2000.06.18 14:44	Indian Ocean	So	-13.80	97.45	14	7.5	15	29	7.9	7.8	8.1	39	7.8		7.8	
2000.10.06 04:30	W Honshu	So	35.46	133.13	10	6.5	15	12	6.7	6.8	6.8	53	6.9		6.9	
2000.11.16 04:54	New Ireland	I	-3.98	152.16	33	7.6	24	80	8.0	7.5	7.7	136	7.8	Yes	8.0	-5.7
2000.11.17 21:01	New Britain	I	-5.50	151.78	37	7.4	17	47	7.8	7.5	7.7	76	7.6	Yes	7.7	-5.3
2001.01.26 03:16	S India (Bhuj)	R	23.42	70.23	10	7.6	20	28	7.6	7.8	8.0	31	7.5		7.5	-4.5
2001.02.28 18:54	Washington	P	47.14	-122.72	52	N/A	51	9	6.8	6.6	6.7	8	6.7		6.7	
2001.03.24 06:27	W Honshu	P	34.08	132.53	49	6.7	47	17	6.8	7.0	7.0	25	6.9		6.9	
2001.06.23 20:33	Peru	I	-16.27	-73.64	8	8.3	30	138	8.4	7.5	7.7	156	8.0	Yes	8.4	-5.3
2002.08.19 11:08	Fiji Islands	D	-23.88	178.50	676	7.6	699	21	7.7	7.5	7.7	13	7.6		7.6	-3.1
2002.11.03 22:12	Alaska	RS	63.52	-147.44	4	N/A	15	94	7.9	7.4	7.6	31	7.4		7.4	
2003.01.22 02:06	Mexico	I	18.84	-103.82	24	7.6	26	29	7.5	7.5	7.6	28	7.4	Yes	7.5	-4.3
2003.05.21 18:44	N Algeria	R	36.96	3.63	9	6.7	15	20	6.8	7.0	7.1	23	6.8		6.8	-5.2
2003.07.15 20:27	Carlsberg Ridge	So	-2.56	68.30	10	N/A	15	94	7.5	7.4	7.5	102	7.5		7.5	-6.0
2003.08.04 04:37	Scotia Sea	No	-60.56	-43.49	10	7.1	15	45	7.6	7.3	7.5	45	7.4		7.4	-5.1
2003.09.25 19:50	Hokkaido	I	41.82	143.91	13	8.1	28	64	8.3	7.9	8.2	82	7.9	Yes	8.3	-4.6
2003.09.27 11:33	Siberia	S	50.04	87.81	1	7.3	15	22	7.2	7.4	7.5	77	7.3		7.3	
2003.11.17 06:43	Rat Islands	I	51.15	178.65	5	7.7	22	48	7.7	7.4	7.5	73	7.5	Yes	7.7	-5.3
2003.12.26 01:56	S Iran	S	29.00	58.31	10	6.5	15	11	6.6	6.7	6.7	23	6.6		6.6	
2004.11.11 21:26	Timor	I	-8.17	124.86	10	7.4	17	34	7.5	7.3	7.4	52	7.4	Yes	7.5	-5.1
2004.11.26 02:25	Papua Indonesia	P	-3.57	135.35	10	6.9	12	18	7.1	7.0	7.1	25	7.2		7.2	
2004.11.28 18:32	Hokkaido	I	43.00	145.06	39	7.0	47	10	7.0	7.2	7.3	18	7.1		7.1	-4.3
2004.12.23 14:59	Macquarie	So	-49.31	161.35	35	7.9	28	53	8.1	7.8	8.1	64	7.9		7.9	-4.9
2004.12.26 00:58	Sumatra-Andaman	IT?	3.30	95.98	39	8.2	29	278	9.3	8.1	8.3	418	8.6	Yes	9.2	-5.4
2005.02.05 12:23	Celebes Sea	D	5.36	123.21	501	7.0	531	9	7.1	N/A	N/A	15	7.0		7.0	
2005.03.02 10:42	Banda Sea	W	-6.53	129.94	201	7.1	196	9	7.1	7.0	7.1	11	7.1		7.1	
2005.03.28 16:09	N Sumatra	I	2.09	97.11	21	8.1	30	110	8.6	8.2	8.6	108	8.2	Yes	8.6	-4.4
2005.06.13 22:44	Chile	W	-19.99	-69.20	115	7.8	95	13	7.7	7.6	7.8	18	7.7		7.7	
2005.06.15 02:50	N California	So	41.284	-125.983	10	7.1	20	24	7.2	6.9	7.0	34	7.2		7.2	-5.1
2005.07.24 15:42	Nicobar	So	7.92	92.19	16	7.1	12	20	7.2	7.2	7.3	33	7.3		7.3	
2005.08.16 02:46	Honshu	I	38.28	142.04	36	7.0	37	24	7.2	7.4	7.5	54	7.3	Yes	7.3	
2005.09.09 07:26	New Ireland	W	-4.54	153.45	91	7.4	84	58	7.6	7.5	7.7	144	7.7		7.7	-6.3
2005.09.26 01:55	N Peru	W	-5.67	-76.41	127	7.5	108	13	7.5	7.5	7.6	21	7.5		7.5	
2005.10.08 03:50	Pakistan	R	34.54	73.59	26	7.3	12	21	7.6	7.6	7.8	57	7.4		7.4	
2005.11.14 21:38	E Honshu	P	38.10	144.93	11	6.8	18	16	7.0	7.1	7.2	19	7.1		7.1	
2006.01.02 06:10	S Sandwich Islands	So	-60.81	-21.47	10	7.1	20	28	7.4	7.2	7.3	38	7.4		7.4	
2006.01.27 16:58	Banda Sea	D	-5.48	128.09	397	7.5	397	22	7.6	7.5	7.7	21	7.6		7.6	
2006.02.22 22:19	Mozambique	N	-21.32	33.58	11	7.0	12	14	7.0	7.3	7.5	20	7.0		7.0	
2006.04.20 23:25	Koryakia	Ro	61.08	167.09	22	7.3	12	31	7.6	7.3	7.4	38	7.4		7.4	-4.9
2006.05.03 15:26	Tonga	W	-20.13	-174.16	55	7.9	68	47	8.0	7.7	7.9	44	7.9		7.9	-4.3
2006.05.16 10:39	Kermadec	D	-31.78	-179.31	151	7.4	155	26	7.4	7.5	7.6	27	7.5		7.5	
2006.07.17 08:19	Indonesia	T	-9.25	107.41	34	7.2	20	139	7.7	7.2	7.3	178	7.5	Yes	7.7	-6.5
2006.08.20 03:41	Scotia Sea	So	-61.01	-34.39	10	7.0	17	18	7.0	6.9	7.0	17	7.0		7.0	-4.5
2006.09.28 06:22	Samoa Islands	P	-16.57	-172.04	39	6.7	12	11	6.9	7.0	7.1	13	6.9		6.9	-4.3
2006.11.15 11:14	Kuril Islands	I	46.68	153.22	28	7.9	13	106	8.3	7.6	7.7	123	7.9	Yes	8.2	-5.2
2006.12.26 12:26	Taiwan	P	21.83	120.54	10	7.1	23	16	6.9	6.9	7.0	19	7.0		7.0	
2006.12.26 12:34	Taiwan	P	22.01	120.51	10	N/A	34	17	6.8	7.0	7.1	34	7.2		7.2	
2007.01.13 04:23	Kuril Islands	P	46.29	154.45	10	7.9	12	56	8.1	7.8	8.1	88	8.0		8.0	-5.1
2007.01.21 11:27	Molucca Sea	P?	1.24	126.40	22	7.3	22	39	7.5	7.4	7.5	37	7.4		7.4	-4.9
2007.01.30 04:54	Macquarie	So	-54.89	145.73	10	6.8	14	13	6.8	6.7	6.7	17	6.8		6.8	-4.7
2007.04.01 20:39	Solomon Islands	I	-8.45	156.96	10	N/A	23	89	8.1	7.5	7.7	114	7.9	Yes	8.2	-5.1
2007.08.01 17:08	Vanuatu	W	-15.74	167.75	120	7.2	127	19	7.2	7.0	7.0	81	7.4		7.4	-5.9
2007.08.08 17:04	Java	W	-5.97	107.66	289	7.4	304	29	7.5	7.4	7.5	21	7.4		7.4.	



**Table 2**

Magnitude corrections for event depth (PREM / PREM\_NC)

Depth range (km)	Correction (magnitude units)
< 15 continental crust	-0.28 (not used in this study)
< 24.4 continental crust	-0.15
< 24.4 other types	no correction
24.4-220	no correction
220-271	+0.05
271-371	+0.06
371-400	+0.07
400-471	+0.12
471-571	+0.15
571-671	+0.18
$\geq 671$	+0.22

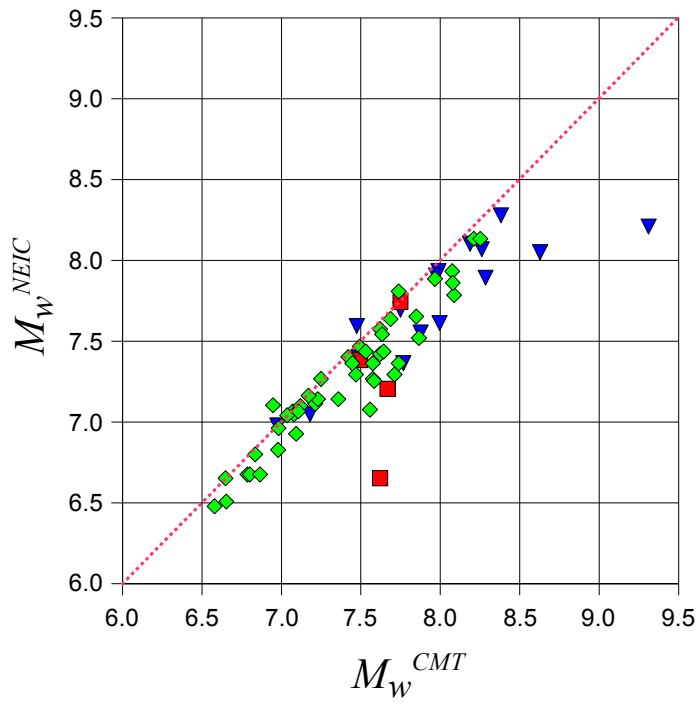


Figure 1 a)

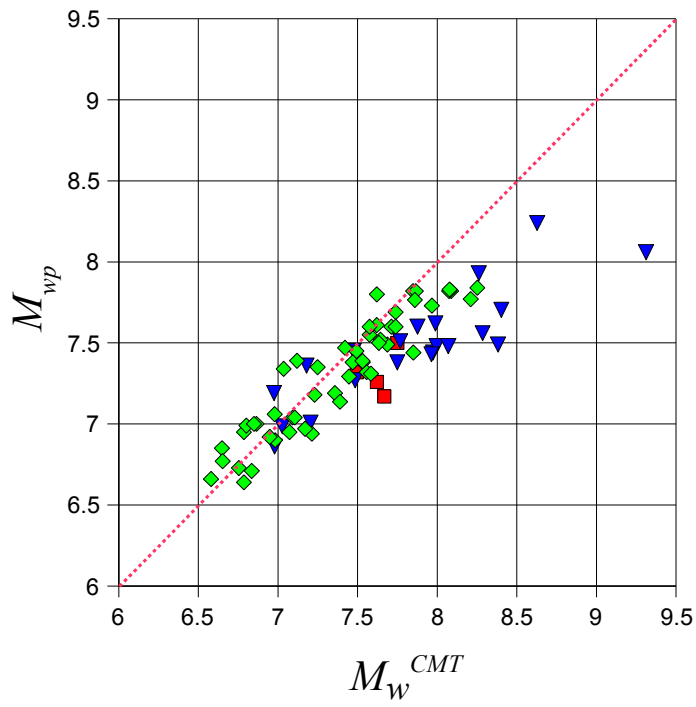


Figure 1 b)

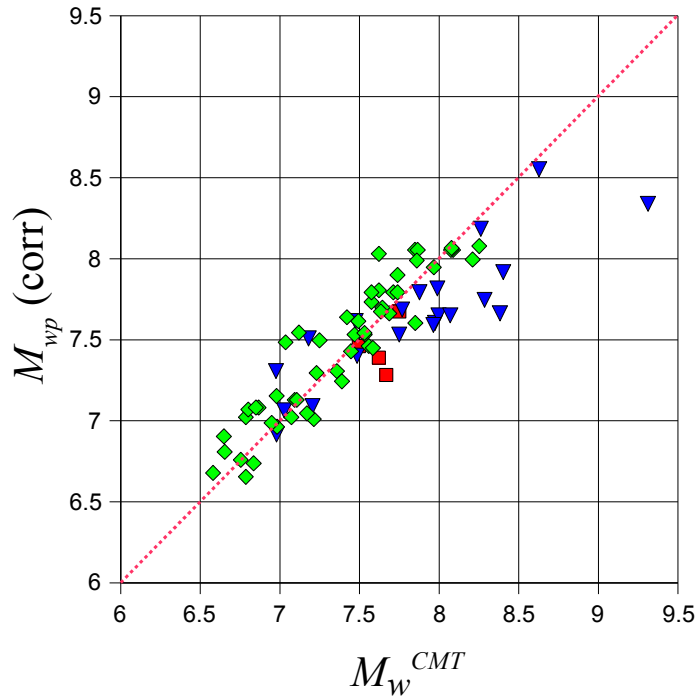
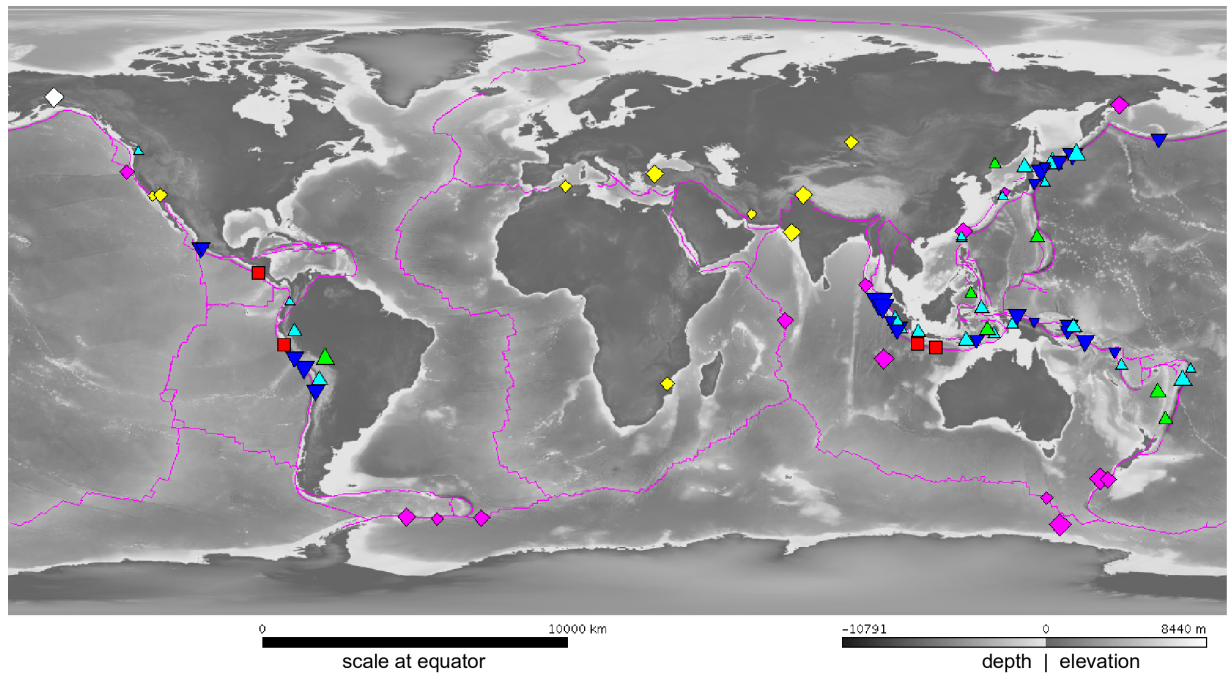


Figure 1 c)

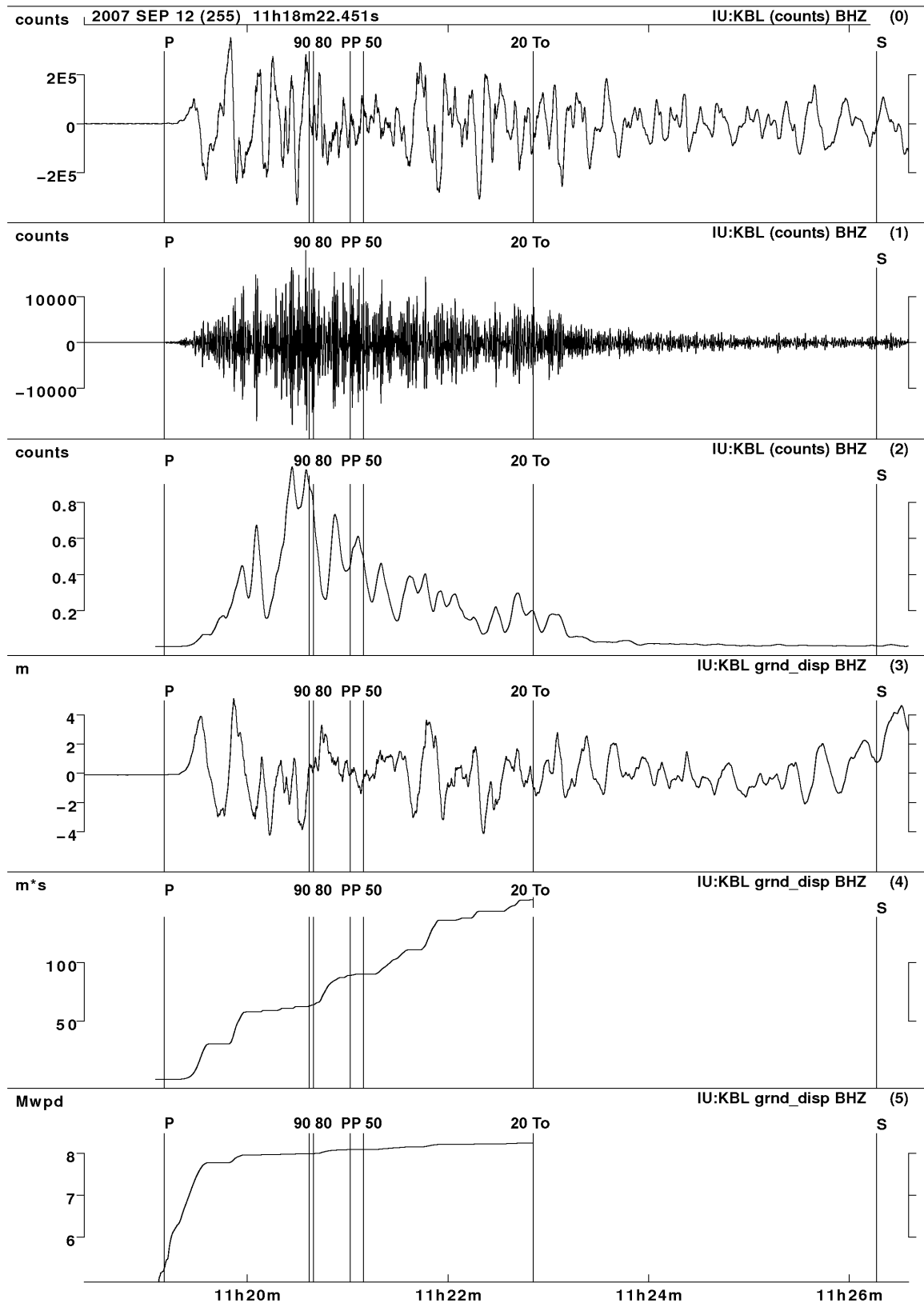
## Figure 1

Moment magnitudes from rapid analysis methods using seismic  $P$ -waves compared to CMT magnitude  $M_w^{CMT}$  for the studied events (Table 1, Figure 2). a)  $M_w^{NEIC}$  from the NEIC Fast Moment Tensor procedure (Sipkin, 1994; <http://earthquake.usgs.gov>); b)  $M_{wp}$  from this study, determined following the procedure described by Tsuboi (2000), Hirshorn (2006) and Lomax *et al.* (2007); c)  $M_{wp}$  from this study with magnitude dependent correction of Whitmore *et al.* (2002). Event symbols are: interplate thrust events (blue inverted triangles); tsunami earthquakes (red squares); other event types (green diamonds). In this and the following figures the value  $M_w^{CMT}=9.3$  for 2004.12.26 Sumatra-Andaman is from Tsai *et al.* (2005).



**Figure 2**

World map showing earthquakes used in this study (*c.f.* Table 1). Symbols show earthquake type: I - interplate thrust (blue inverted triangles); T - tsunami earthquake (red squares); W – down-dip and P – intraplate (light blue triangles); D – deep (green triangles); So - strike-slip oceanic, Ro – reverse-faulting oceanic and No – normal-faulting oceanic (magenta diamonds); S - strike-slip continental, R - reverse-faulting continental and N - normal-faulting continental (yellow diamonds); hybrid events (white diamonds). Base map from NGDC (2006); plate boundaries (magenta lines) from Coffin *et al.* (1998).



**Figure 3**

Duration-amplitude processing steps for the 12 September 2007, M8.4 Sumatra earthquake recorded at station IU:KBL at 49° GCD to the northwest of the event. Trace (0): raw, velocity seismogram; Trace (1): 1.0 Hz, Gaussian-filtered seismogram; Trace (2): smoothed, velocity-squared envelope; Trace (3): amplitude-corrected, ground-displacement seismogram; Trace (4): integral of trace (3) over the source duration using Equation (3) before multiplication by  $k$  and  $C_M$ , note that for this seismogram the integral over positive values of displacement,  $u^+(t)$ , in trace (3) gives the maximum result; Trace (5): Raw  $M_{pwd}$  magnitude obtained from trace (3) using Equation (2).  $P$ ,  $PP$  and  $S$  indicate the PREM\_NC predicted arrival times for the first arriving,  $P$ ,  $PP$  and  $S$  waves from the hypocentre. 90, 80, 50 and 20 indicate the times at which the envelope function, trace (2), last drops below 90% ( $T^{90}$ ), 80% ( $T^{80}$ ), 50% ( $T^{50}$ ) and 20% ( $T^{20}$ ) of its peak value, respectively;  $T_0$  indicates the estimated apparent duration,  $T_0$ , for this station. See Appendix B for more details. The  $PP$  amplitudes on this recording (visible around 11h21m to 11h22m) are larger relative to the  $P$  amplitudes than they are for most other recordings for this or other events.

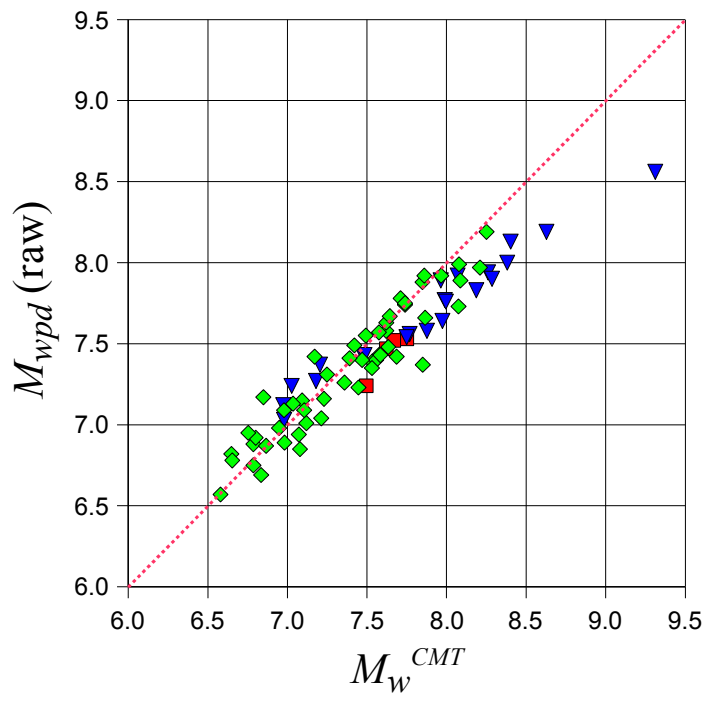


Figure 4 a)

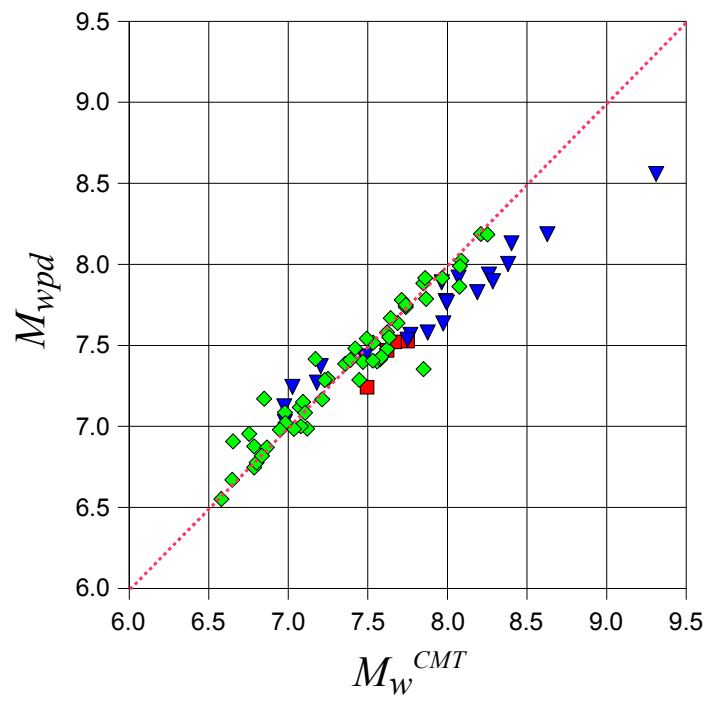
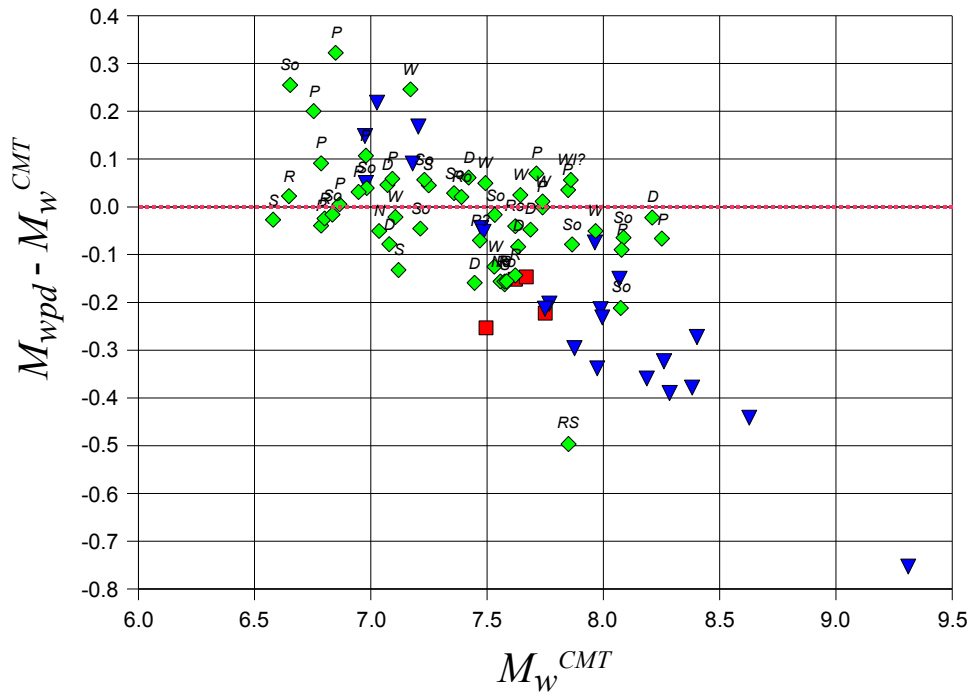


Figure 4 b)



4 c)

## Figure 4

Results for duration-amplitude magnitude  $M_{wpd}$  with no moment scaling for interplate thrust or tsunami events (*i.e.*, application of Equation 3) for the studied events (Table 1). a) “raw”  $M_{wpd}$  given by direct application of Equation (3) without any corrections for event type compared to CMT magnitude  $M_w^{CMT}$ ; b)  $M_{wpd}$  (with event type corrections) compared to  $M_w^{CMT}$ ; c)  $\Delta M = M_{wpd} - M_w^{CMT}$  compared to  $M_w^{CMT}$ ;  $\Delta M$  has a standard deviation of  $\sigma = 0.17$ . Material properties at the source are corrected to correspond to the PREM or PREM\_NC model values at the CMT centroid depth (Table 2). The comparison between  $M_0^{pd}$  and  $M_0^{CMT}$  to determine  $k$  in Equation (3) excludes interplate thrust and tsunami events and 2002.11.03 Alaska (labelled RS in plots) which has a poor  $T_0$  estimate due to exceptional source complexity (*e.g.*, Fuis and Wald, 2003). Event symbols and labels as in Figures 1 and 2.



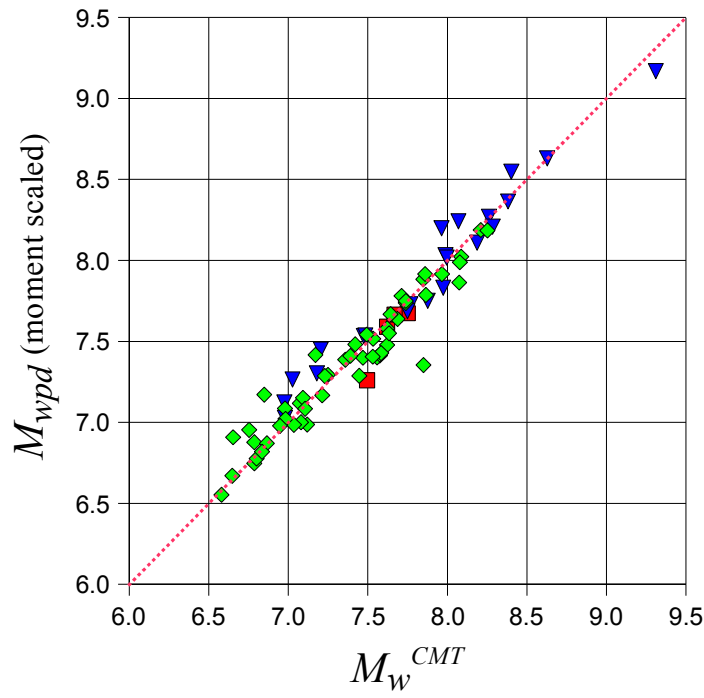
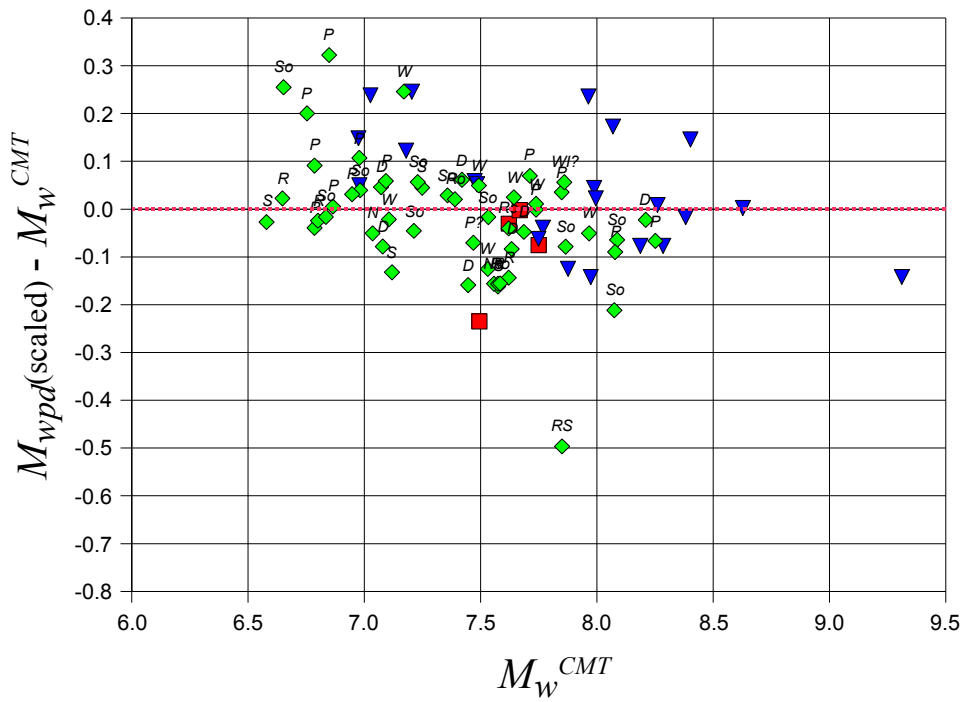


Figure 5 a)

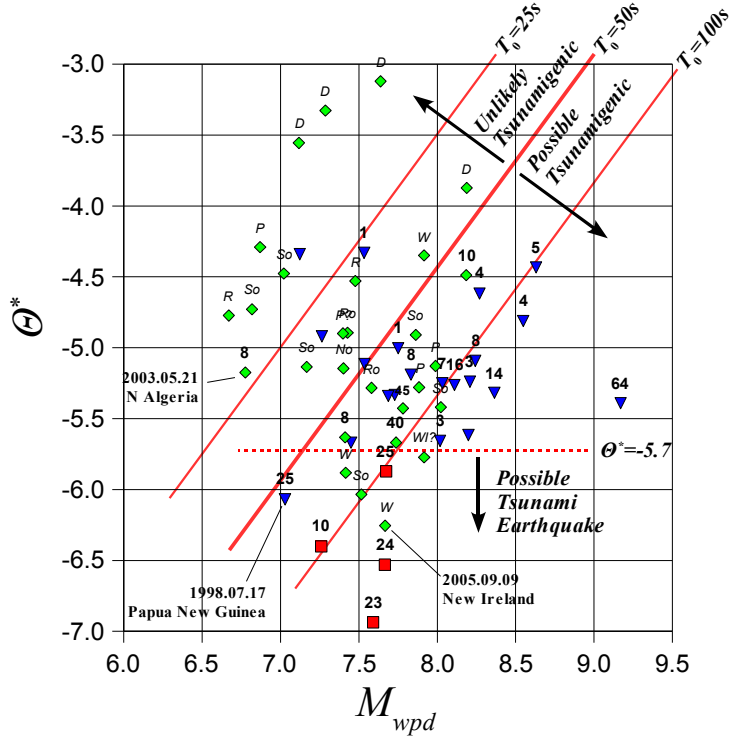


5 b)

## Figure 5

Results for duration-amplitude magnitude  $M_{wpd}$  corrected with moment scaling for interplate thrust and tsunami events (*i.e.*, application of Equations 3 and 5a or 5b) for the studied events

(Table 1). a)  $M_{wpd}$  compared to CMT magnitude  $M_w^{CMT}$ , b)  $\Delta M = M_{wpd} - M_w^{CMT}$  compared to  $M_w^{CMT}$ .  $\Delta M$  has a standard deviation of  $\sigma = 0.11$ . Material properties at the source are corrected to correspond to the PREM or PREM\_NC model values at the CMT centroid depth (Table 2). Event symbols and labels as in Figures 1 and 2.



**Figure 6**

$\Theta^*$  values from application of Equation (8) to duration-amplitude results with moment scaling for interplate thrust and tsunami events (*i.e.*, application of Equations 3 and 5a or 5b) for the studied events where  $\sigma_{T_0} < 2T_0/3$  or  $I_t \geq 1$  (Table 1).  $\Theta^*$  values are plotted against  $M_{wpd}$ . Bold numbers show the measure of tsunami importance,  $I_t$ , based on maximum water height and descriptive indices of tsunami effects from the NOAA/WDC Historical Tsunami Database (NGDC, 2008). Event symbols as in Figure 1 and event type labels (shown for events with  $I_t < 1$ ) as in Figure 2 and Table 1. Dashed red line shows  $\Theta^* = -5.7$  cutoff for identification of tsunami earthquakes; lines constant  $T_0$  are shown in red; thick red line shows  $T_0 = 50$  sec cutoff for identification of tsunamigenic earthquakes.

## Appendix A - Far-field estimation of seismic moment from $P$ waveforms

Following Aki and Richards (1980), Boatwright and Choy (1986), Tsuboi *et al.* (1995) and Kanamori and Rivera (2004), if  $u(t)$  is the amplitude corrected, far-field,  $P$ -displacement for an earthquake source of duration  $T_0$ , then a theoretical expression for scalar seismic moment,  $M_0$ , is,

$$M_0 = C_M \int_{t_p}^{t_p + T_0} u(t) dt. \quad (A1)$$

In the above expression  $t_p$  is the  $P$  arrival time and  $u(t)$  is corrected for geometrical spreading and attenuation.  $C_M = 4\pi\rho_s^{1/2}\rho_r^{1/2}\alpha_s^{5/2}\alpha_r^{1/2}Ff_s$ , where  $\rho$  and  $\alpha$  are the density and  $P$  wave speed, respectively, at the source  $s$  or the recording station  $r$ , and  $F$  and  $f_s$  are corrections for radiation pattern and free-surface amplification, respectively.

In this study we use the 1-D, spherical, PREM model (Dziewonski and Anderson, 1981) without a crust (PREM\_NC) for amplitude correction of the displacement waveforms for attenuation and geometrical spreading. We choose PREM so that we can make unbiased comparisons with the Global CMT results, which are based on PREM. In PREM\_NC the crustal layers are replaced by a layer with the PREM properties of the uppermost mantle; this eliminates unrealistic, discontinuous jumps in material properties and magnitude estimates at the crustal boundaries for small changes in the nominal hypocenter depth. Calculations are initially performed using  $\rho$  and  $\alpha$  values for the uppermost mantle; later, for shallow continental events and deeper events (Table 1), the effect of the crust or event depth on  $\rho$  and  $\alpha$  is re-introduced as a magnitude correction using the PREM properties at the depth of the event (Table 2).

The geometrical spreading is calculated from the spreading of rays between the source and station in the PREM\_NC model using a standard expression (*e.g.*, Aki and Richards, 1980, eq. 9.44; Shearer, 1999, eq. 6.23).

The attenuation correction is made in the frequency domain using standard relations (*e.g.*, Shearer, 1999; Lay, 2002),

$$A_{corr}(\omega) = A_0(\omega) e^{-\omega t^*/2}, \quad (A2)$$

and,

$$t^* = \int_{path} \frac{dt}{Q(r)}, \quad (A3)$$

where  $A_0(\omega)$  and  $A_{corr}(\omega)$  are the Fourier transforms of the initial and attenuation corrected displacements, and the integral in Equation (A3) is taken using the source-station ray path and corresponding  $Q$  values from the PREM\_NC model.

If the integral in Equation A1 includes all of the  $P$  wave group ( $P$ ,  $pP$  and  $sP$ ) then the correction to displacement for radiation pattern is given by a factor  $F = \sqrt{[\langle(F^P)^2\rangle]/(F^{gP})^2}$  where  $\langle(F^P)^2\rangle = 4/15$  (*e.g.*, Boatwright and Choy, 1986) is the mean square radiation coefficient for  $P$  waves, and  $F^{gP}$  is a generalized radiation pattern coefficient for the  $P$  wave group. For observations at teleseismic distances Newman and Okal (1998) suggest a constant value  $F^{gP}=1$  for the generalized radiation coefficient which is appropriate for dip-slip faulting but

considered too high by as much as a factor of 4 for strike-slip faulting (Boatwright and Choy, 1986; Choy and Boatwright, 1995). This choice of  $F^{gP}$  gives  $F = \sqrt{[\langle (F^P)^2 \rangle]} = \sqrt{(4/15)} \approx 0.52$ . However, if the integral in Equation A1 includes only the direct  $P$  waves, then  $F = \sqrt{[1/\langle (F^P)^2 \rangle]} = \sqrt{(15/4)} \approx 1.9$  (e.g., Tusboi *et al.*, 1999). Since in this study we compensate for the presence of non direct  $P$  waves by taking the integral in Equation (1) separately over the positive and negative displacement amplitudes, we use  $F = \sqrt{(15/4)}$  for the radiation pattern correction for non-strike slip events. Because of the ambiguity noted above in the radiation coefficients for strike-slip faulting, we determine empirically a magnitude correction for strike-slip events so that their  $M_{wpd}$  magnitudes best match  $M_w^{CMT}$  on average (see appendix B).

The correction for free-surface amplification at the station site introduces an additional factor of  $f_s = 1/2$ . Incorporating the corrections for radiation pattern and free-surface amplification in  $C_M$ , and using PREM\_NC upper-mantle material properties for the source,  $\rho_s = 3.38 \text{ g/cm}^3$ ,  $\alpha_s = 8.10 \text{ km/sec}$ , and PREM upper crust properties for the recording stations,  $\rho_r = 2.60 \text{ g/cm}^3$ ,  $\alpha_r = 5.80 \text{ km/sec}$ , gives  $C_M = 1.62 \times 10^{19}$  when geometrical spreading is expressed as an equivalent source-station distance in units of km. The magnitude corrections to account for  $\rho_s$  and  $\alpha_s$  for shallow continental and deeper sources are listed in Table 2.

In the preceding we have not directly accounted for the  $PP$  phase which arrives in the  $P$ -wave group and may be expected to bias the moment estimates upwards. However, the effect of  $PP$  on the duration-amplitude magnitude calculations seems to be minor to insignificant, for three main reasons: 1) The majority of large events used in our study (69 out of 79) have duration  $T_0 < 2 \text{ min}$ . For these events,  $PP$  at stations with  $GCD > \sim 50^\circ$  arrives later than the window  $t_p$  to  $t_p + T_0$  used for integration in Equation A1. Thus  $PP$  is for the most part not included in the calculation. Two of the few events where  $PP$  may be included in the integral are the 2004.12.26 M9.3 Sumatra-Andaman ( $T_0 \approx 400\text{-}500 \text{ sec}$ ) and the 2006.07.17 M7.7 Indonesia tsunami earthquake ( $T_0 \approx 180 \text{ sec}$ ); but for both of these events raw  $M_{wpd}$  is less than  $M_w^{CMT}$  and the moment corrected  $M_{wpd} \approx M_w^{CMT}$ , so there is no evidence of overestimation of  $M_w$  due to neglecting the effects of  $PP$ . 2) An examination of the displacement signals for longer duration ( $T_0 > 2 \text{ min}$ ) and larger events shows that the  $PP$  amplitudes are always smaller than the  $P$  amplitudes (e.g., Figure 3, trace (3) exhibits a relatively large  $PP$  signal) and, for a majority of traces, are so small as to be difficult to identify. This phenomenon may be due to destructive interference of  $PP$  pulses for longer duration events, since  $PP$  is related to the Hilbert transform of the  $P$  waveform. The Hilbert transform of a simple pulse is a quasi-symmetric pair of positive and negative pulses. For the  $PP$  case, these two pulses are separated by a time interval that is much less than  $T_0$  for the studied events, consequently there can be cancellation between positive and negative  $PP$  pulses originating from different sub-sources of the evolving rupture. 3) The moment magnitude is related to the logarithm of the moment (cf., Equation 2; Figure 3, traces 4 and 5), so a relatively large error in moment corresponds to a small error in magnitude (e.g., a factor of two error in  $M_0$  leads to a change in  $M_w$  of 0.2).

## Appendix B – Duration-amplitude moment and magnitude calculation

For each earthquake we assume that we have a hypocentre location and predicted  $P$  and  $S$  travel times from the hypocentre to each recording station. Currently, most real-time monitoring agencies have this information within a few minutes after OT for local and regional events (GCD to stations  $< \sim 30^\circ$ ), and within about 10 to 15 minutes of OT for teleseismic events (GCD to stations  $> \sim 30^\circ$ ). We also assume that we have available vertical-component, broadband, digital seismograms for about 20 or more stations at  $30^\circ$  to  $90^\circ$  GCD from the source, and that these stations are moderately well distributed in distance and azimuth to avoid biases due to rupture directivity and other effects. We exclude from the analysis poor quality seismograms that are noisy, clipped, truncated, or otherwise corrupted.

For the present study we examine a set of recent earthquakes with a large range of magnitudes ( $M_w^{CMT}$  6.6 to 9.3) and diverse source types (Table 1). For each event, we obtain from the IRIS Data Management Center a set of broadband vertical (BHZ) component recordings at stations from  $30^\circ$  to  $90^\circ$  GCD from the event. Typically we use about 20 to 50 records, selecting stations well distributed in distance for events which have more than 50 available records. All averages and standard deviations are obtained using robust statistics (*i.e.*, 20% trimmed - rejection of the upper and lower 20 percentiles of values), typically data from 15 to 45 stations are retained.

### Duration determination

At teleseismic distance, direct  $P$ -waves contain much more, higher-frequency energy than do other wave types such as  $pP$ ,  $sP$ ,  $PP$  or  $S$ . In consequence, the duration of the direct  $P$  waves and an apparent source duration,  $T_0$ , can be obtained from high-frequency seismograms (Lomax, 2005; Lomax *et al.*, 2007). We exploit this behaviour to estimate  $T_0$  for each station using vertical-component seismograms, with the following procedure (see also Figure 3), based on that of Lomax (2005) and Lomax *et al.* (2007): 1) Convert the seismograms from each station to high-frequency records using a narrow-band, Gaussian filter of the form  $e^{-\alpha(|f-f_{cent}|/f)^2}$ , where  $f$  is frequency,  $f_{cent}$  the filter center frequency, and  $\alpha$  sets the filter width. Here we use  $f_{cent} = 1.0$  Hz and  $\alpha = 10.0$ ; as in Lomax (2005) and Lomax *et al.* (2007), in contrast to the 2-4 Hz band-pass filter used by Hara (2007). 2) Convert the high-frequency seismogram to velocity-squared time-series by squaring each of the data values. 3) Smooth the velocity-squared time-series with a 10 sec wide, triangle function to form a station envelope function. 4) Measure the set of time delays after the  $P$  time at which the envelope function last drops below 90% ( $T^{90}$ ), 80% ( $T^{80}$ ), 50% ( $T^{50}$ ) and 20% ( $T^{20}$ ) of its peak value. 5) Calculate the apparent source duration,  $T_0$ , for the station using the following algorithm,

$$T_0 = (1 - w) T^{90} + w T^{20}, \quad (B1)$$

where the weight  $w = [(T^{80} + T^{50}) / 2 - 20 \text{ sec}] / 40 \text{ sec}$ , with limiting values  $0 \leq w \leq 1$ .

The form of  $w$  and choice of 20% and 90% of the envelope peak value to measure  $T_0$  follow from examination of the shape of the summary envelope functions used in this study. In general, the 20% peak value gives better agreement with published results for the larger events (*e.g.*,  $T_0 > 100$  sec), while the 90% peak value better results for the smallest events (*e.g.*,  $T_0 < 100$  sec).

100 sec), in comparison to twice the CMT centroid minus origin times and other estimates of source duration. The necessity for different treatment of smaller and larger events is due to the longer length of the exponentially decaying,  $P$  coda in proportion to the source duration for smaller events than for larger events.

We also calculate an average  $T_0$  and associated standard deviation for each event by taking the geometric mean and geometric standard deviation of the station  $T_0$  estimates using robust statistics (*i.e.*, 20% trimmed measures).

In general the duration-amplitude  $T_0$  estimates are greater than twice the CMT centroid minus origin times (*cf.* Table 1), since the  $T_0$  estimation procedure accounts well for the very long and complex rupture of larger events (which are not well represented by the single-triangle source function used in CMT), while  $T_0$  will tend to overestimate the durations for smaller events due to problems with the relatively long coda in the high-frequency seismograms.

### ***Duration-amplitude moment and magnitude calculation***

We evaluate the seismic moment,  $M_0^{pd}$ , for each station using vertical-component seismograms and the following procedure (see also Figure 3): 1) Bandpass from 1 to 200 sec (see Appendix C), remove the instrument response and apply geometrical spreading and attenuation corrections to convert each seismogram to amplitude corrected, ground displacement. 2) Cut each seismogram from 10 seconds before the  $P$ -arrival to the  $P$ -arrival time plus the source duration,  $T_0$ , or to 10 seconds before the  $S$  arrival, whichever is earlier, to obtain  $P$ -wave seismograms. 3) Apply Equations (3) and (5a or 5b) to each  $P$ -wave seismograms to obtain station moment estimates. 4) Multiply the station moment value by a factor  $T_0 / t_{S-P}$  if  $T_0 > t_{S-P}$ , where  $t_{S-P}$  is the  $S$  arrival time minus the  $P$  arrival time. We calculate an average  $M_0^{pd}$  and associated standard deviation for each event by taking the geometric mean and geometric standard deviation of the station moment estimates using robust statistics (*i.e.*, 20% trimmed - rejection of the upper and lower 20 percentiles of values). We calculate the duration-amplitude magnitude,  $M_{wpd}$ , through application of the standard moment to moment magnitude relation (Hanks and Kanamori, 1979; Bormann, 2002),

$$M_{wpd} = (\log_{10} M_0^{pd} - 9.1) / 1.5, \quad (B2)$$

where  $M_0^{pd}$  has units of N-m.

We include a constant,  $k$ , in Equation (3) to compensate for the errors and biases in the geometrical spreading and attenuation corrections and in the terms of  $C_M$ . We evaluate  $k$  through comparison of our  $M_0^{pd}$  values for each event against the corresponding CMT moment values,  $M_0^{CMT}$ , so that the mean of  $\log_{10}(M_0^{pd}/M_0^{CMT}) \rightarrow 0$ , giving  $k \approx 1.2$ . This evaluation excludes interplate thrust, tsunami and strike-slip events and 3 November 2002 Alaska (labelled RS in plots) which has an unstable  $T_0$  estimate due to exceptional source complexity (*e.g.*, Fuis and Wald, 2003). We use only interplate thrust and tsunami events to determine the constant  $M_0^{cutoff}$  in Equation (5a), giving  $M_0^{cutoff} \approx 7.5 \times 10^{19}$  N-m (equivalent to  $M_w \approx 7.2$ ), and to determine the optimal value of  $R$  in Equation (4) by minimizing the standard deviation of  $\log_{10}(M_0^{pd}/M_0^{CMT})$ , giving  $R \approx 0.45$ . The optimal values of  $M_0^{cutoff}$  and  $R$  are sensitive to details of the algorithms used to estimate  $T_0$  and moment; a change of  $\pm 0.25$  in  $R$  gives about half the variance reduction relative to  $R = 0$  (*i.e.*, no moment scaling) than gives  $R \approx 0.45$ . The empirically determined magnitude correction to account for the radiation pattern of strike-slip events (types S and So in Table 1) has a value of 0.13 magnitude units; this value

implies that for strike-slip events an additional factor of about 1.6 is needed in the correction for radiation pattern,  $F$ , in Equation (A1).

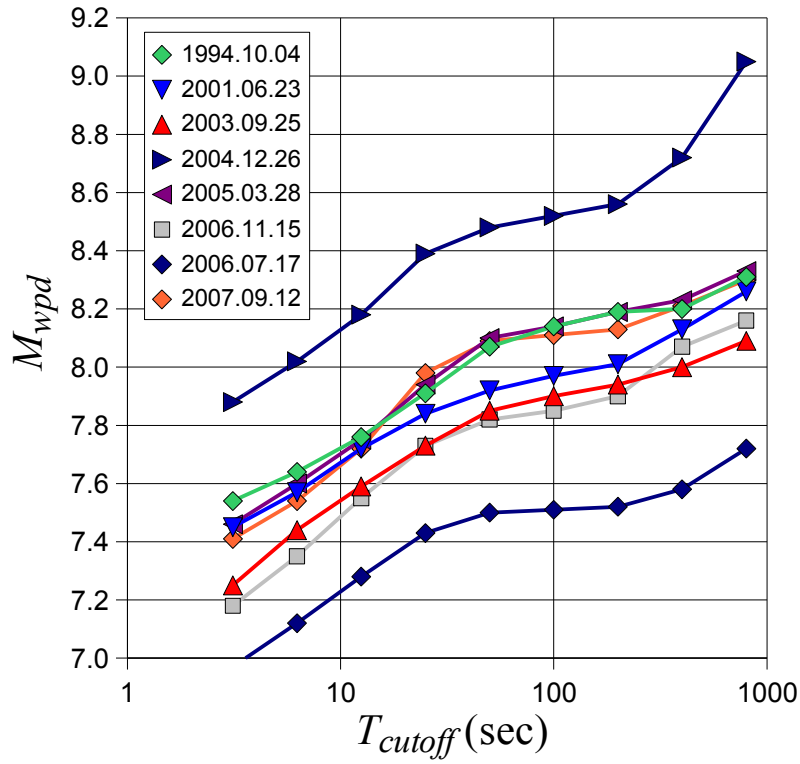


## Appendix C – Dependence of duration-amplitude results on long-period cutoff

The values of moment and of moment magnitude,  $M_{wpd}$ , for large events obtained with the duration-amplitude procedure depend on the long-period cutoff used when processing the seismograms. Indeed, it is generally accepted that magnitude saturation, regardless of the magnitude estimation technique, is related to the long-period, data cutoff being lower than a corner period above which the displacement spectrum flattens to an amplitude proportional to the static moment (*e.g.*, Stein and Okal, 2006). Magnitude saturation also arises for methods that use a signal duration after the initial  $P$  arrival that is shorter than the duration of significant  $P$  signal and the source duration (*e.g.*, Granville *et al.*, 2005); the duration-amplitude procedure avoid this latter problem by explicitly taking into account the source duration.

Figure C1 shows duration-amplitude magnitudes,  $M_{wpd}$ , with no moment scaling, for the 7 largest and one tsunami earthquake from the studied events, plotted as a function of long-period cutoff,  $T_{cutoff}$ . With increasing  $T_{cutoff}$  to about 50 sec there is an increase in magnitude estimates for all events; this increase can be associated with magnitude saturation due to  $T_{cutoff}$  being lower than the long-period spectral corner for  $P$  waves. However, at around  $T_{cutoff} = 50$ -200 sec the curves in Figure C1 flatten and the magnitude estimates are nearly independent of  $T_{cutoff}$ , indicating that the long-period corner for  $P$  waves for each event has been reached and that the resulting magnitude estimates should not be saturated. Above around  $T_{cutoff} = 200$  sec the magnitude estimates again increase with  $T_{cutoff}$ , examination of the processed seismograms shows that this increase is primarily an artefact of amplification of long-period noise in the  $P$ -wave train during the removal of the instrument response. The onset of  $P$ -wave noise above about 200 sec period is expected since the typical long-period corner is about 120 to 360 sec for the very-broadband instruments providing much of the data used in this study.

These results and Figure C1 indicate that: 1) The optimal long-period cutoff for the studied data set is 100-200 sec. 2) The trend of increasing underestimate of  $M_w^{CMT}$  by unscaled  $M_{wpd}$  with increasing  $M_w^{CMT}$  (Figure 4) cannot be attributed to a magnitude saturation problem due to insufficient, long-period signal.



**Figure C1**

Duration-amplitude magnitudes,  $M_{wpd}$ , with no moment scaling (*i.e.*, application of Equation 3) for the 7 largest and one tsunami earthquake (2006.07.17 Indonesia) from the studied events (Table 1) plotted as a function of long period cutoff used for analysis. The events are identified by their origin dates.